

Spherical Harmonics

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OST FHO Campus Rapperswil

Spring Semester 2022

Spherical Harmonics

Seminar der Ostschweizer Fachhochschule in Rapperswil, Frühjahrsemester 2022 dem Thema Spezielle Funktionen war, die grosse Vielfalt von speziellen Funktionsfamilien zu ergründen, die im Laufe der Zeit für Anwendungen erdacht wurden. Dieses Buch bringt esungsteils mit den von den Seminarteilnehmern anarbeiten zusammen.

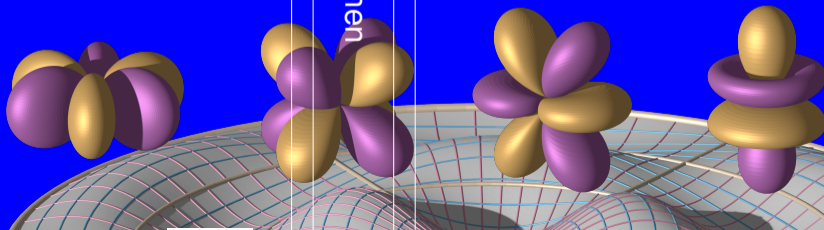
Spezielle Funktionen

Mathematisches Seminar

Spezielle Funktionen

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Spherical Harmonics *and* Electron Orbitals

Mathematisches Seminar

hydrogen $n=5, l=2, m=1$


MINILITE
PHYSICS

Seminar der Ostschweizer Fachhochschule im Frühjahrsemester 2022 dem Thema "Spherical Harmonics" war, die grosse Vielfalt von Familien zu ergründen, die in Anwendungen erdacht wurden, esungsteils mit den von den Mitarbeiterinnen zusammen.

Wave
function
 $\psi_{nlm}(x,y,z)$

possible
electron
positions

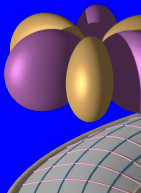


Table of Contents

- 1 Fourier on \mathbb{R}^2
- 2 The functions $Y_{m,n}(\varphi, \vartheta)$
- 3 Fourier on S^2
- 4 Quantum Mechanics

Definition

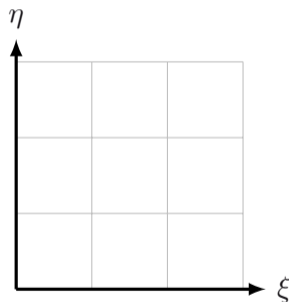
A function

$$f : \mathbb{R}^2 \rightarrow \mathbb{C}$$

is a “nice periodic function” when it is

- smooth,
- differentiable,
- (abs.) integrable,
- periodic on $[0, 1] \times [0, 1]$, i.e.

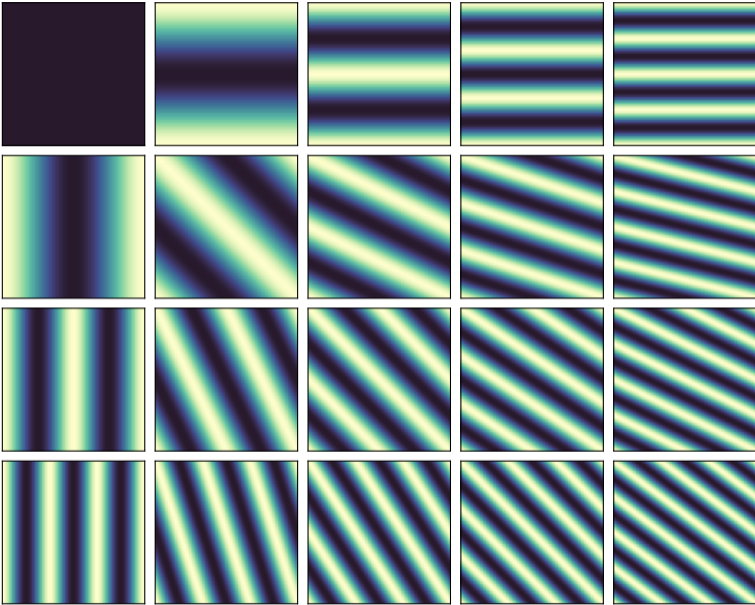
$$f(\xi, \eta) = f(\xi + 1, \eta) = f(\xi, \eta + 1).$$



Basis Functions

The space of nice periodic functions is spanned by the (also nice) functions

$$B_{m,n}(\xi, \eta) = e^{i2\pi m\xi} e^{i2\pi n\eta}.$$



Definition

Let $f(\xi, \eta)$ and $g(\xi, \eta)$ be nice periodic functions. Their inner product is

$$\langle f, g \rangle = \iint_{[0,1]^2} f(\xi, \eta) \bar{g}(\xi, \eta) d\xi d\eta.$$

Inner Product

Definition

Let $f(\xi, \eta)$ and $g(\xi, \eta)$ be nice periodic functions. Their inner product is

$$\langle f, g \rangle = \iint_{[0,1]^2} f(\xi, \eta) \bar{g}(\xi, \eta) d\xi d\eta.$$

Definition

For a nice periodic function $f(\xi, \eta)$: the numbers

$$c_{m,n} = \langle f, B_{m,n} \rangle$$

are the *Fourier coefficients* or *spectrum* of f .

Theorem

For nice periodic functions:

$$f(\xi, \eta) = \sum_{m \in \mathbb{Z}} \sum_{n \in \mathbb{Z}} c_{m,n} B_{m,n}(\xi, \eta)$$

where

$$c_{m,n} = \langle f, B_{m,n} \rangle.$$

Why exponentials?

Why $B_{m,n} = e^{i2\pi m\xi} e^{i2\pi n\eta}$?

Because ∇^2

Fourier's Problem

$$\nabla^2 f(\xi, \eta) = \frac{\partial^2 f}{\partial \xi^2} + \frac{\partial^2 f}{\partial \eta^2} = \lambda f(\xi, \eta)$$

Fourier's Problem

$$\nabla^2 f(\xi, \eta) = \frac{\partial^2 f}{\partial \xi^2} + \frac{\partial^2 f}{\partial \eta^2} = \lambda f(\xi, \eta)$$

Solution

Separation ansatz:

$$f(\xi, \eta) = M(\xi)N(\eta)$$

Resulting ODEs:

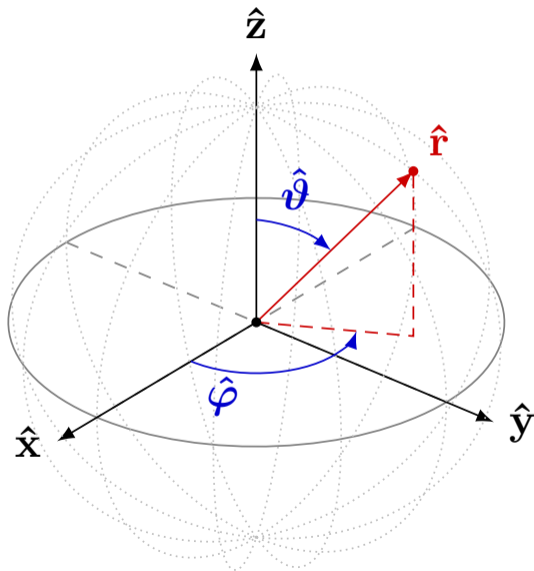
$$\frac{d^2 M}{d\xi^2} = \kappa M(\xi),$$

$$\frac{d^2 N}{d\eta^2} = (\lambda - \kappa)N(\eta)$$

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Spherical Coordinates



Variables

$$r \in \mathbb{R}^+$$

$$\vartheta \in [0, \pi]$$

$$\varphi \in [0, 2\pi)$$

To cartesian

$$x = r \cos \varphi \sin \vartheta$$

$$y = r \sin \varphi \sin \vartheta$$

$$z = r \cos \vartheta$$

Spherical Laplacian

Cartesian Laplacian

$$\nabla^2 := \frac{\partial^2}{\partial \xi^2} + \frac{\partial^2}{\partial \eta^2}$$

Spherical Laplacian

Cartesian Laplacian

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Spherical Laplacian

$$\nabla^2 := \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2} \left[\frac{1}{\sin \vartheta} \frac{\partial}{\partial \vartheta} \left(\sin \vartheta \frac{\partial}{\partial \vartheta} \right) + \frac{1}{\sin^2 \vartheta} \frac{\partial^2}{\partial \varphi^2} \right]$$

Spherical Laplacian

Cartesian Laplacian

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Spherical Laplacian

$$\nabla^2 := \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2} \underbrace{\left[\frac{1}{\sin \vartheta} \frac{\partial}{\partial \vartheta} \left(\sin \vartheta \frac{\partial}{\partial \vartheta} \right) + \frac{1}{\sin^2 \vartheta} \frac{\partial^2}{\partial \varphi^2} \right]}_{\text{Surface Spherical Laplacian } \nabla_s^2}$$

Spherical Laplacian

Cartesian Laplacian

$$\nabla^2 := \frac{\partial^2}{\partial \xi^2} + \frac{\partial^2}{\partial \eta^2}$$

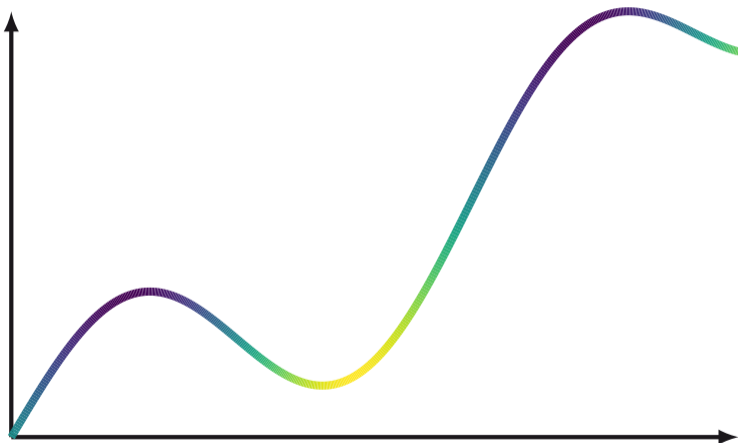
Spherical Laplacian

$$\nabla^2 := \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \underbrace{\frac{1}{r^2} \left[\frac{1}{\sin \vartheta} \frac{\partial}{\partial \vartheta} \left(\sin \vartheta \frac{\partial}{\partial \vartheta} \right) + \frac{1}{\sin^2 \vartheta} \frac{\partial^2}{\partial \varphi^2} \right]}_{\text{Surface Spherical Laplacian } \nabla_s^2}$$

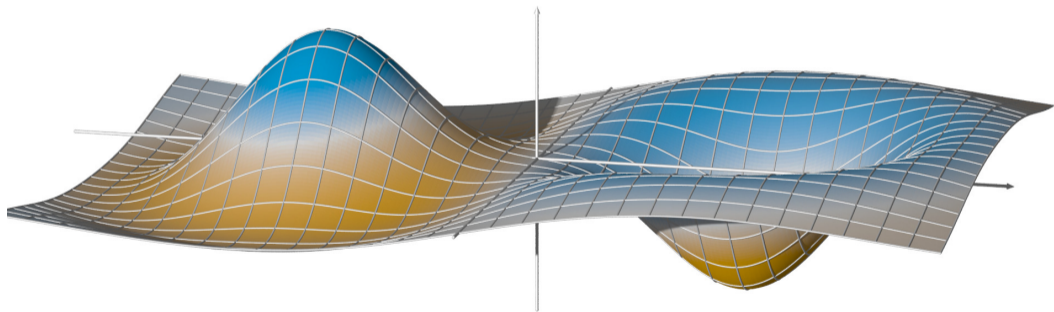
Surface Spherical Laplacian

$$\nabla_s^2 := r^2 \nabla^2 - \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right)$$

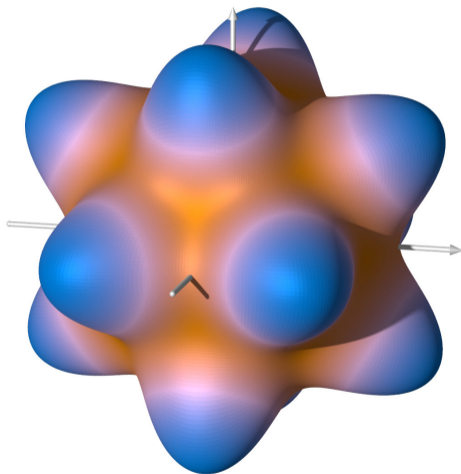
Geometrical Intuition



Geometrical Intuition

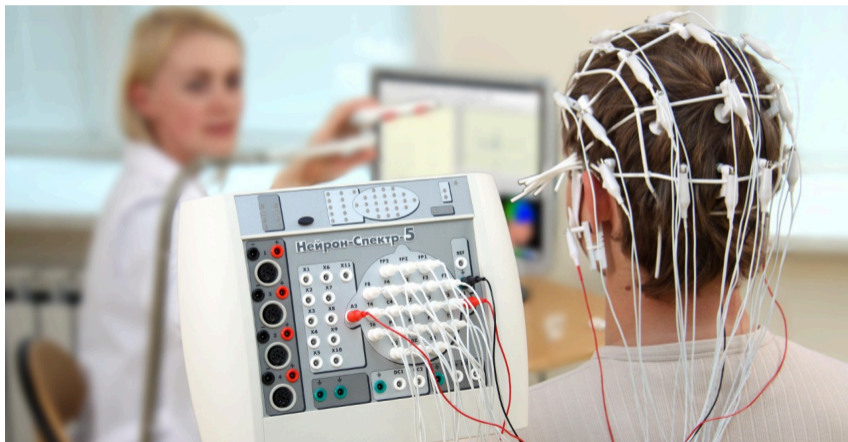


Geometrical Intuition



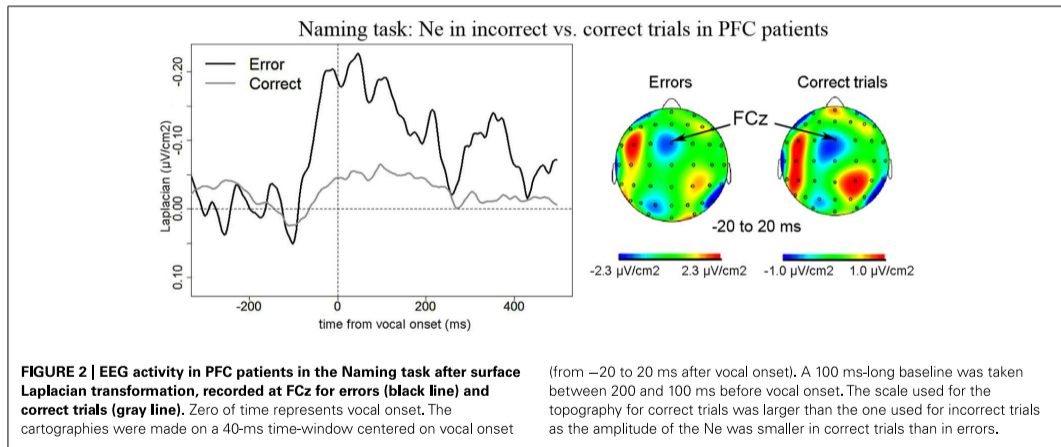
Where is ∇_s^2 useful?

To do brain scans, apparently [2]



Where is ∇_s^2 useful?

To do brain scans, apparently [2]

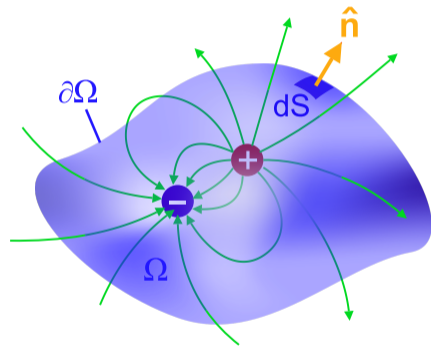


Electrodynamics

$$\nabla^2 \phi = \nabla \cdot \nabla \phi \quad \left(\phi = \int_A^B \mathbf{E} \cdot d\mathbf{l} \right)$$

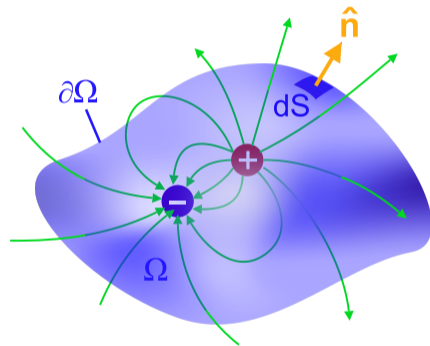
Electrodynamics

$$\begin{aligned}\nabla^2\phi &= \nabla \cdot \nabla\phi \\ &= \nabla \cdot \mathbf{E}\end{aligned}\quad \left(\phi = \int_A^B \mathbf{E} \cdot d\mathbf{l}\right)$$



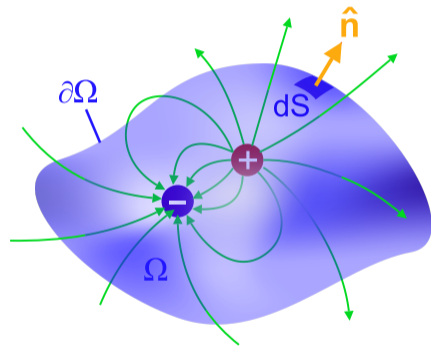
Electrodynamics

$$\begin{aligned}\nabla^2\phi &= \nabla \cdot \nabla\phi && \left(\phi = \int_A^B \mathbf{E} \cdot d\mathbf{l} \right) \\ &= \nabla \cdot \mathbf{E} \\ &= \int_{\Omega} (\nabla \cdot \mathbf{E}) \cdot d\mathbf{s} = \oint_{\partial\Omega} \mathbf{E} \cdot d\mathbf{s}\end{aligned}$$



Electrodynamics

$$\begin{aligned}\nabla^2 \phi &= \nabla \cdot \nabla \phi && \left(\phi = \int_A^B \mathbf{E} \cdot d\mathbf{l} \right) \\ &= \nabla \cdot \mathbf{E} \\ &= \int_{\Omega} (\nabla \cdot \mathbf{E}) \cdot d\mathbf{s} = \oint_{\partial\Omega} \mathbf{E} \cdot d\mathbf{s} \\ &= \frac{\rho}{\epsilon}\end{aligned}$$



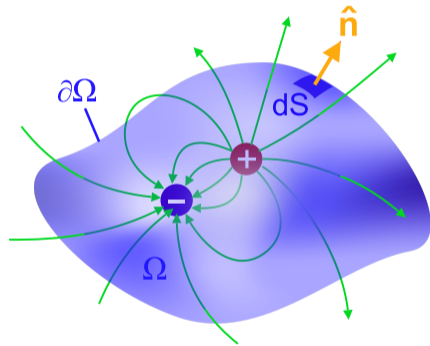
Brain Scans

Electrodynamics

$$\begin{aligned}\nabla^2\phi &= \nabla \cdot \nabla\phi && \left(\phi = \int_A^B \mathbf{E} \cdot d\mathbf{l} \right) \\ &= \nabla \cdot \mathbf{E} \\ &= \int_{\Omega} (\nabla \cdot \mathbf{E}) \cdot d\mathbf{s} = \oint_{\partial\Omega} \mathbf{E} \cdot d\mathbf{s} \\ &= \frac{\rho}{\epsilon}\end{aligned}$$

So over the scalp

$$\nabla_s^2\phi = \frac{\rho_s}{\epsilon} = \text{Current flow in the brain}$$



New Hard Problem

The Problem

$$\nabla_s^2 f(\varphi, \vartheta) = \lambda f(\varphi, \vartheta)$$

The Problem

$$\frac{1}{\sin \vartheta} \frac{\partial}{\partial \vartheta} \left(\sin \vartheta \frac{\partial f}{\partial \vartheta} \right) + \frac{1}{\sin^2 \vartheta} \frac{\partial^2 f}{\partial \varphi^2} = \lambda f(\varphi, \vartheta)$$

New Hard Problem

The Problem

$$\frac{1}{\sin \vartheta} \frac{\partial}{\partial \vartheta} \left(\sin \vartheta \frac{\partial f}{\partial \vartheta} \right) + \frac{1}{\sin^2 \vartheta} \frac{\partial^2 f}{\partial \varphi^2} = \lambda f(\varphi, \vartheta)$$

Idea

Separation ansatz:

$$f(\varphi, \vartheta) = \Phi(\varphi)\Theta(\vartheta)$$

From the “easy” part:

$$\frac{d^2 \Phi}{d\varphi^2} = \kappa \Phi(\varphi) \implies \Phi(\varphi) = e^{im\varphi}, \quad m \in \mathbb{Z}$$

Associated Legendre Differential Equation

Separation (cont.)

The hard part is the ODE for $\Theta(\vartheta)$:

$$\sin^2 \vartheta \frac{d^2 \Theta}{d(\cos \vartheta)^2} - 2 \cos \vartheta \frac{d\Theta}{d \cos \vartheta} + \left[n(n+1) - \frac{m^2}{\sin^2 \vartheta} \right] \Theta(\cos \vartheta) = 0$$

Associated Legendre Differential Equation

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Substituting $x = \cos \vartheta$ and $y = \Theta$:

Definition (Associated Legendre Differential Equation)

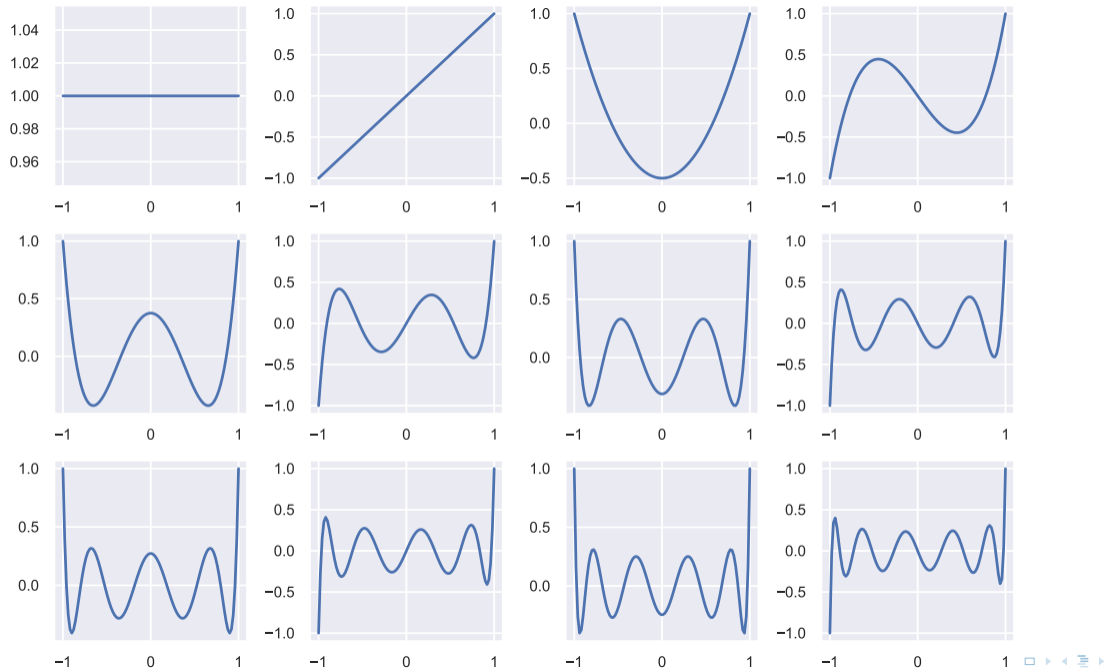
$$(1-x^2) \frac{d^2 y}{dx^2} - 2x \frac{dy}{dx} + \left[n(n+1) - \frac{m^2}{1-x^2} \right] y(x) = 0$$

Definition (Legendre Polynomials)

The polynomials

$$\begin{aligned} P_n(x) &= \sum_{k=0}^{\lfloor n/2 \rfloor} \frac{(-1)^k (2n-2k)!}{2^n k! (n-k)! (n-2k)!} x^{n-2k} \\ &= {}_2F_1 \left(\begin{matrix} n+1, & -n \\ 1 & \end{matrix}; \frac{1-x}{2} \right) \\ &= \frac{1}{n! 2^n} \frac{d^n}{dx^n} (x^2 - 1)^n \end{aligned}$$

are a solution to the associated Legendre differential equation when $m = 0$.



Associated Legendre Polynomials

Lemma

For $x \in [-1, 1]$ the polynomials

$$P_{m,n}(x) = (1 - x^2)^{m/2} \frac{d^m}{dx^m} P_n(x)$$

solve the associated Legendre differential equation.

Associated Legendre Polynomials

Lemma

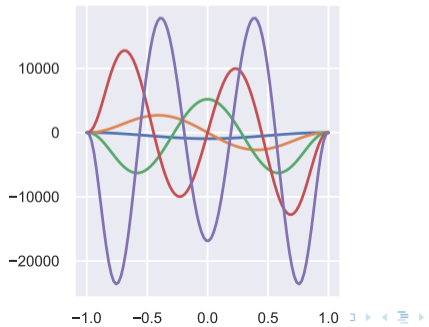
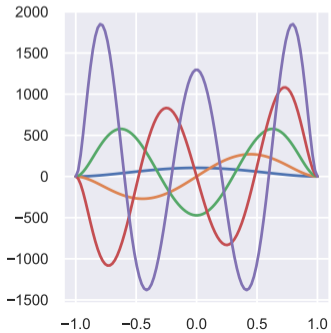
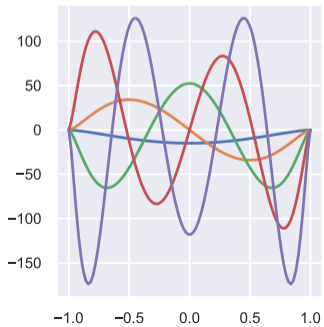
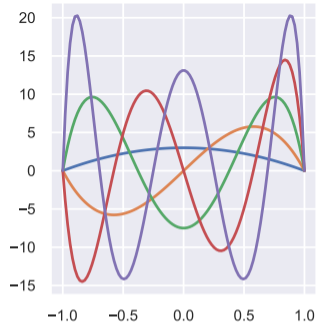
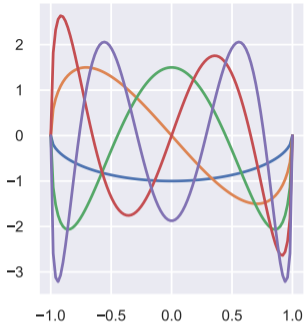
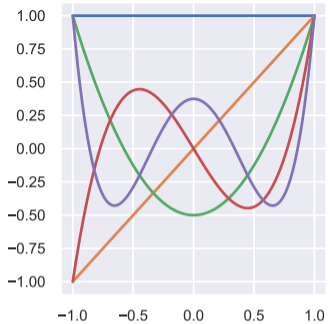
For $x \in [-1, 1]$ the polynomials

$$P_{m,n}(x) = (1 - x^2)^{m/2} \frac{d^m}{dx^m} P_n(x)$$

solve the associated Legendre differential equation.

Observation

If $m > n$ then $P_{m,n}(x) = 0$ for all x .



The Problem

$$\nabla_s^2 f(\varphi, \vartheta) = \lambda f(\varphi, \vartheta)$$

Current solution

For $m \in \mathbb{Z}$ and $m < n$:

$$\tilde{Y}_{m,n}(\varphi, \vartheta) = \Phi(\varphi)\Theta(\vartheta) = e^{im\varphi} P_{m,n}(\cos \vartheta)$$

What do they look like?

Python Magic

Intuition of conditions for m and n

Recurrence Relation(s)?

$$\tilde{Y}_{m+1,n} \stackrel{?}{=} f(\tilde{Y}_{m,n}, \tilde{Y}_{m-1,n}, \tilde{Y}_{m,n-1}, \dots)$$

$$\tilde{Y}_{m,n+1} \stackrel{?}{=} f(\tilde{Y}_{m,n}, \tilde{Y}_{m-1,n}, \tilde{Y}_{m,n-1}, \dots)$$

$$\tilde{Y}_{m+1,n+1} \stackrel{?}{=} f(\tilde{Y}_{m,n}, \tilde{Y}_{m-1,n}, \tilde{Y}_{m,n-1}, \dots)$$

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Basis functions?

The functions $\tilde{Y}_{m,n}$ span the space of nice functions $S^2 \rightarrow \mathbb{C}$.

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Definition

The inner product of nice functions $f(\varphi, \vartheta)$ and $g(\varphi, \vartheta)$ from S^2 to \mathbb{C} is

$$\langle f, g \rangle = \iint_{S^2} f(\varphi, \vartheta) \bar{g}(\varphi, \vartheta) d\Omega$$

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$$\langle f, g \rangle = \iint_{S^2} f(\varphi, \vartheta) \bar{g}(\varphi, \vartheta) d\Omega = \int_0^{2\pi} \int_0^{\pi} f(\varphi, \vartheta) \bar{g}(\varphi, \vartheta) \sin \vartheta d\vartheta d\varphi$$

Orthonormality

Definition

A set of basis functions are *orthonormal* if

$$\langle B_{m,n}, B_{m',n'} \rangle = \begin{cases} 1 & m = m' \wedge n = n' \\ 0 & \text{else} \end{cases}$$

Orthonormality

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Problem

$$\langle \tilde{Y}_{m,n}, \tilde{Y}_{m',n'} \rangle = \begin{cases} \frac{4\pi}{2n+1} \frac{(n+m)!}{(n-m)!} & m = m' \wedge n = n' \\ 0 & \text{else} \end{cases}$$

Spherical Harmonics

Definition

The orthonormal spherical harmonics are

$$Y_{m,n}(\varphi, \vartheta) = N_{m,n} e^{im\varphi} P_{m,n}(\cos \vartheta)$$

where the normalisation constant

$$N_{m,n} = \sqrt{\frac{2n+1}{4\pi} \frac{(n-m)!}{(n+m)!}}$$

Fixed

$$\langle Y_{m,n}, Y_{m',n'} \rangle = \begin{cases} 1 & m = m' \wedge n = n' \\ 0 & \text{else} \end{cases}$$

Theorem

For nice periodic functions on S^2 :

$$f(\varphi, \vartheta) = \sum_{m \in \mathbb{Z}} \sum_{n \in \mathbb{Z}} c_{m,n} Y_{m,n}(\varphi, \vartheta)$$

where

$$c_{m,n} = \langle f, Y_{m,n} \rangle.$$

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Linear and Rotational Kinetic Energy

Momentum and KE

$$\mathbf{p} = m\mathbf{v}, \quad E_k = \frac{\mathbf{p}^2}{2m}$$

Linear and Rotational Kinetic Energy

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Angular Momentum and KE

$$\mathbf{L} = \mathbf{r} \times \mathbf{p}, \quad E_{k,a} = \frac{\mathbf{L}^2}{2mr^2}$$

Linear and Rotational Kinetic Energy

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QM Formulation

$$\hat{\mathbf{p}} = -i\hbar\nabla, \quad \hat{E}_k = -\frac{\hbar^2}{2m}\nabla^2$$

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Angular Momentum and KE

$$\mathbf{L} = \mathbf{r} \times \mathbf{p}, \quad E_{k,a} = \frac{\mathbf{L}^2}{2mr^2}$$

QM Formulation

Pretty long derivation yields:

$$\hat{E}_{k,a} = -\frac{\hbar^2}{2mr^2}\nabla_s^2$$

Intuition for the Operators

Schrödinger Equation

Time independent SE

$$\hat{\mathcal{H}}|\Psi\rangle = E|\Psi\rangle$$

Schrödinger Equation

Time independent SE

$$\left(\hat{E}_k + U\right) |\Psi\rangle = E|\Psi\rangle$$

Schrödinger Equation

Time independent SE

$$\left(\frac{\hat{\mathbf{p}}^2}{2m} + U \right) |\Psi\rangle = E|\Psi\rangle$$

Schrödinger Equation

Time independent SE

$$\text{Meili} \quad \left[-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + U(x) \right] \Psi(x) = E\Psi(x)$$

Schrödinger Equation

Time independent SE

$$3D \quad \left[-\frac{\hbar^2}{2m} \nabla^2 + U(\mathbf{x}) \right] \Psi(\mathbf{x}) = E \Psi(\mathbf{x})$$

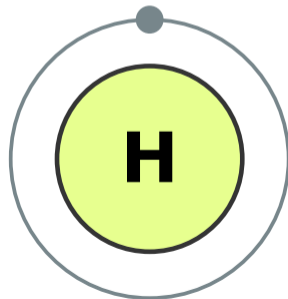
Schrödinger Equation

Time independent SE

$$\left\{ -\frac{\hbar^2}{2m} \frac{1}{r^2} \left[\nabla_s^2 - \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) \right] + U(\mathbf{r}) \right\} \Psi(\mathbf{r}) = E\Psi(\mathbf{r})$$

But why?

Hydrogen atom has radial symmetry!



Schrödinger Equation

Time independent SE

$$\left[\frac{\hat{\mathbf{L}}^2}{2mr^2} + \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + U(\mathbf{r}) \right] \Psi(\mathbf{r}) = E\Psi(\mathbf{r})$$

Time independent SE

$$\left[\hat{E}_{k,a} + \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + U(\mathbf{r}) \right] \Psi(\mathbf{r}) = E \Psi(\mathbf{r})$$

Schrödinger Equation

Time independent SE

$$\left[\underbrace{\hat{E}_{k,a} + \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right)}_{\text{Kinetic Energy}} + U(\mathbf{r}) \right] \Psi(\mathbf{r}) = E \Psi(\mathbf{r})$$

Time independent SE

$$\left[\hat{E}_{k,a} + \underbrace{\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right)}_{\text{Radial KE } \hat{E}_{k,r}} + U(\mathbf{r}) \right] \Psi(\mathbf{r}) = E \Psi(\mathbf{r})$$

Schrödinger Equation

Time independent SE

$$\left\{ \hat{E}_{k,a} + \hat{E}_{k,r} + U(\mathbf{r}) \right\} \Psi(\mathbf{r}) = E\Psi(\mathbf{r})$$

Electron Orbitals

Bibliography

- [1] minutephysics, *A better way to picture atoms*, May 19, 2021. [Online]. Available: <https://www.youtube.com/watch?v=W2Xb2GFK2yc> (visited on 05/19/2022).
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