

Spherical Harmonics

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Spherical Harmonics

Seminar der Ostschweizer Fachhochschule in Rapperswil, Frühjahrsemester 2022 dem Thema Spezielle Funktionen. Ziel war, die grosse Vielfalt von speziellen Funktionsfamilien zu ergründen, die im Laufe der Zeit für Anwendungen erdacht wurden. Dieses Buch bringt es zusammen mit den von den Seminarteilnehmern erarbeiteten Beiträgen zusammen.

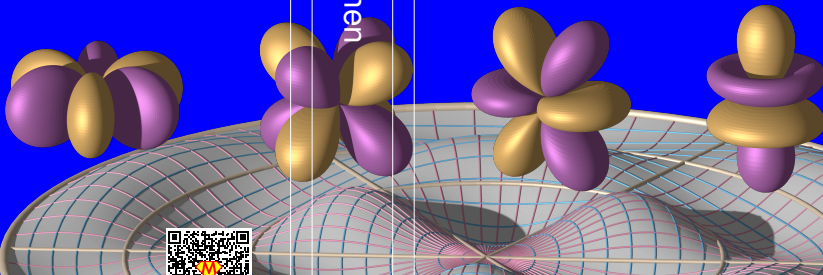
Spezielle Funktionen

Mathematisches Seminar

Spezielle Funktionen

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Spherical Harmonics *and* Electron Orbitals

Mathematisches Seminar

hydrogen



Seminar der Ostschweizer Fachhochschule
Frühjahrssemester 2022 dem Thema
Ziel war, die grosse Vielfalt von
Orbitalfamilien zu ergründen, die in
Anwendungen erdacht wurden.
Lösungsteils mit den von den
Kollegen zusammenarbeiten.

Wave
function
 $\psi_{nlm}(x,y,z)$

possible
electron
positions

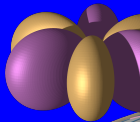


Table of Contents

- 1 Fourier on \mathbb{R}^2
- 2 The functions $Y_{m,n}(\varphi, \vartheta)$
- 3 Fourier on S^2
- 4 Quantum Mechanics

Definition

A function

$$f : \mathbb{R}^2 \rightarrow \mathbb{C}$$

is a “nice periodic function” when it is

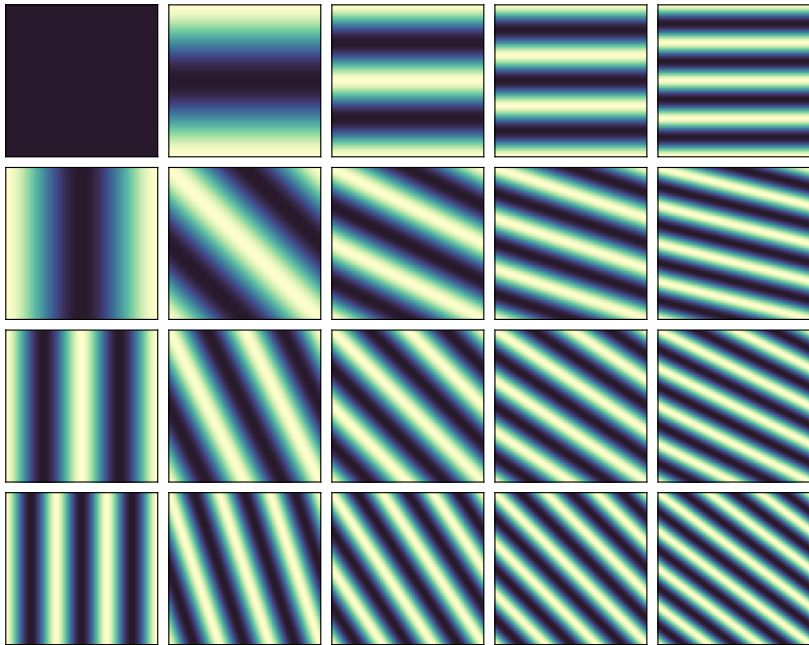
- smooth,
- differentiable,
- (abs.) integrable,
- periodic on $[0, 1] \times [0, 1]$, i.e.

$$f(\mu, \nu) = f(\mu + 1, \nu) = f(\mu, \nu + 1).$$

Basis Functions

The space of nice periodic functions is spanned by the (also nice) functions

$$B_{m,n}(\mu, \nu) = e^{i2\pi m\mu} e^{i2\pi n\nu}.$$



Definition

Let $f(\mu, \nu)$ and $g(\mu, \nu)$ be nice periodic functions. Their inner product is

$$\langle f, g \rangle = \iint_{[0,1]^2} f g^* d\mu d\nu.$$

Inner Product

Definition

Let $f(\mu, \nu)$ and $g(\mu, \nu)$ be nice periodic functions. Their inner product is

$$\langle f, g \rangle = \iint_{[0,1]^2} f g^* d\mu d\nu.$$

Definition

For a nice periodic function $f(\mu, \nu)$: the numbers

$$c_{m,n} = \langle f, B_{m,n} \rangle$$

are the *Fourier coefficients* or *spectrum* of f .

Theorem

For nice periodic functions:

$$f(\mu, \nu) = \sum_{m \in \mathbb{Z}} \sum_{n \in \mathbb{Z}} c_{m,n} B_{m,n}(\mu, \nu)$$

where

$$c_{m,n} = \langle f, B_{m,n} \rangle.$$

Why exponentials?

Why $B_{m,n} = e^{i2\pi m\mu} e^{i2\pi n\nu}$?

Because ∇^2

The Problem

Fourier's Problem

$$\nabla^2 f(\mu, \nu) = \frac{\partial^2 f}{\partial \mu^2} + \frac{\partial^2 f}{\partial \nu^2} = \lambda f(\mu, \nu)$$

The Problem

Fourier's Problem

$$\nabla^2 f(\mu, \nu) = \frac{\partial^2 f}{\partial \mu^2} + \frac{\partial^2 f}{\partial \nu^2} = \lambda f(\mu, \nu)$$

Solution

Separation ansatz:

$$f(\mu, \nu) = M(\mu)N(\nu)$$

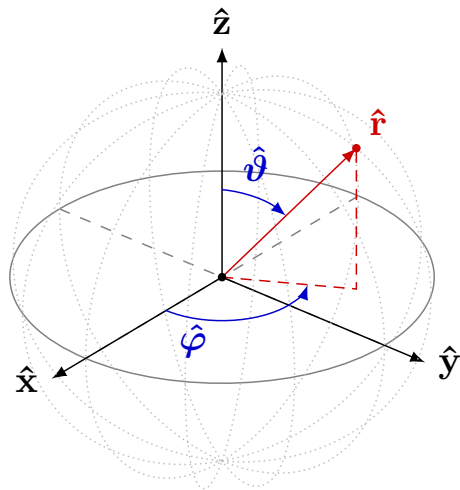
Resulting ODEs:

$$\frac{d^2 M}{d\mu^2} = \kappa M(\mu), \quad \frac{d^2 N}{d\nu^2} = (\lambda - \kappa)N(\nu)$$

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Spherical Coordinates



Variables

$$r \in \mathbb{R}^+$$

$$\vartheta \in [0, \pi]$$

$$\varphi \in [0, 2\pi)$$

To cartesian

$$x = r \cos \varphi \sin \vartheta$$

$$y = r \sin \varphi \sin \vartheta$$

$$z = r \cos \vartheta$$

Spherical Laplacian

Cartesian Laplacian

$$\nabla^2 \equiv \frac{\partial^2}{\partial \mu^2} + \frac{\partial^2}{\partial \nu^2}$$

Spherical Laplacian

Cartesian Laplacian

$$\nabla^2 \equiv \frac{\partial^2}{\partial \mu^2} + \frac{\partial^2}{\partial \nu^2}$$

Spherical Laplacian

$$\nabla^2 \equiv \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2} \left[\frac{1}{\sin \vartheta} \frac{\partial}{\partial \vartheta} \left(\sin \vartheta \frac{\partial}{\partial \vartheta} \right) + \frac{1}{\sin^2 \vartheta} \frac{\partial^2}{\partial \varphi^2} \right]$$

Spherical Laplacian

Cartesian Laplacian

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Spherical Laplacian

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Spherical Laplacian

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Surface Spherical Laplacian

$$\nabla_s^2 \equiv r^2 \nabla^2 - \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right)$$

Geometrical Intuition

Where is ∇_s^2 useful?

To do brain scans, apparently [2]

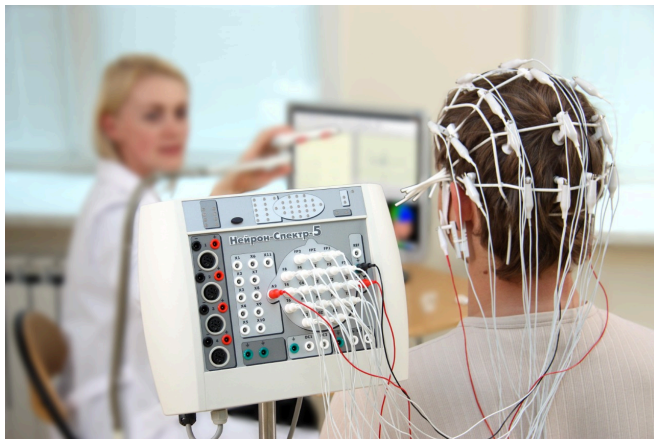


Figure 1: Electroencephalogram (EEG). Image from Wikimedia [3].

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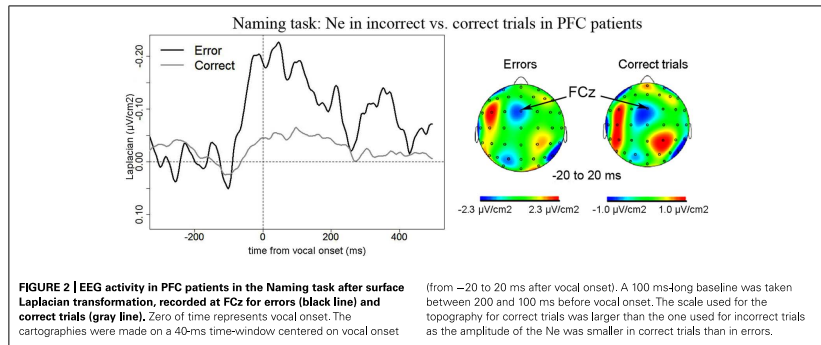


Figure 1: Surface Laplacian in EEG. Taken from [4].

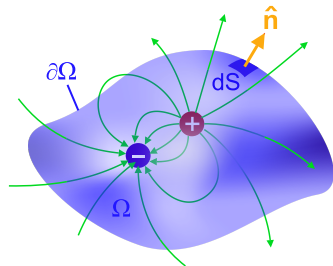
Brain Scans

Electrodynamics

$$\begin{aligned}\nabla^2\phi &= \nabla \cdot \nabla\phi && \left(\phi = \int_A^B \mathbf{E} \cdot d\mathbf{s} \right) \\ &= \nabla \cdot \mathbf{E} \\ &= \int_{\Omega} (\nabla \cdot \mathbf{E}) \cdot d\mathbf{s} = \oint_{\partial\Omega} \mathbf{E} \cdot d\mathbf{s} \\ &= \frac{\rho}{\epsilon}\end{aligned}$$

So over the scalp

$$\nabla_s^2\phi = \frac{\rho_s}{\epsilon} = \text{Current flow in the brain}$$



New Hard Problem

The Problem

$$\nabla_s^2 f(\varphi, \vartheta) = \lambda f(\varphi, \vartheta)$$

New Hard Problem

The Problem

$$\frac{1}{\sin \vartheta} \frac{\partial}{\partial \vartheta} \left(\sin \vartheta \frac{\partial f}{\partial \vartheta} \right) + \frac{1}{\sin^2 \vartheta} \frac{\partial^2 f}{\partial \varphi^2} = \lambda f(\varphi, \vartheta)$$

New Hard Problem

The Problem

$$\frac{1}{\sin \vartheta} \frac{\partial}{\partial \vartheta} \left(\sin \vartheta \frac{\partial f}{\partial \vartheta} \right) + \frac{1}{\sin^2 \vartheta} \frac{\partial^2 f}{\partial \varphi^2} = \lambda f(\varphi, \vartheta)$$

Idea

Separation ansatz:

$$f(\varphi, \vartheta) = \Phi(\varphi)\Theta(\vartheta)$$

From the “easy” part:

$$\frac{d^2 \Phi}{d\varphi^2} = \kappa \Phi(\varphi) \implies \Phi(\varphi) = e^{im\varphi}, \quad m \in \mathbb{Z}$$

Associated Legendre Differential Equation

Separation (cont.)

The hard part is the ODE for $\Theta(\vartheta)$:

$$\sin^2 \vartheta \frac{d^2 \Theta}{d(\cos \vartheta)^2} - 2 \cos \vartheta \frac{d\Theta}{d \cos \vartheta} + \left[n(n+1) - \frac{m^2}{\sin^2 \vartheta} \right] \Theta(\cos \vartheta) = 0$$

Associated Legendre Differential Equation

Separation (cont.)

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Substituting $x = \cos \vartheta$ and $y = \Theta$:

Definition (Associated Legendre Differential Equation)

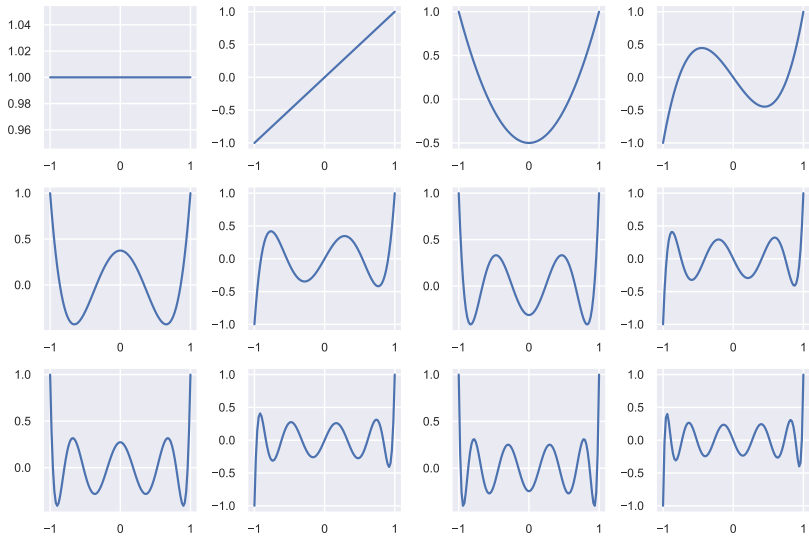
$$(1 - x^2) \frac{d^2 y}{dx^2} - 2x \frac{dy}{dx} + \left[n(n+1) - \frac{m^2}{1 - x^2} \right] y(x) = 0$$

Definition (Legendre Polynomials)

The polynomials

$$\begin{aligned}P_n(x) &= \sum_{k=0}^{\lfloor n/2 \rfloor} \frac{(-1)^k (2n-2k)!}{2^n k! (n-k)! (n-2k)!} x^{n-2k} \\ &= {}_2F_1 \left(\begin{matrix} n+1, & -n \\ 1 & \end{matrix}; \frac{1-x}{2} \right) \\ &= \frac{1}{n! 2^n} \frac{d^n}{dx^n} (x^2 - 1)^n\end{aligned}$$

are a solution to the associated Legendre differential equation when $m = 0$.



Associated Legendre Polynomials

Lemma

For $x \in [-1, 1]$ the polynomials

$$P_{m,n}(x) = (1 - x^2)^{m/2} \frac{d^m}{dx^m} P_n(x)$$

solve the associated Legendre differential equation.

Associated Legendre Polynomials

Lemma

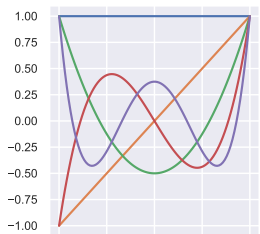
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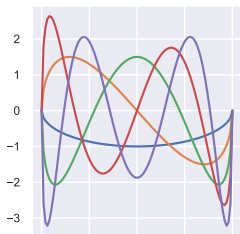
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Observation

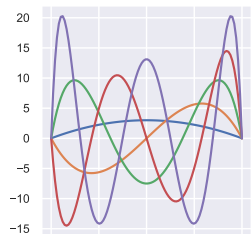
If $m > n$ then $P_{m,n}(x) = 0$ for all x .



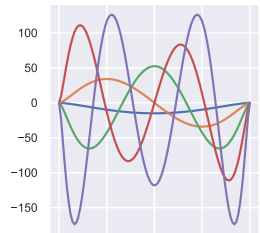
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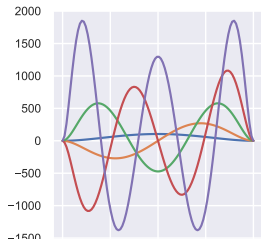
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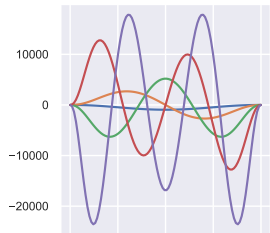
-1.0 -0.5 0.0 0.5 1.0



-1.0 -0.5 0.0 0.5 1.0



-1.0 -0.5 0.0 0.5 1.0



-1.0 -0.5 0.0 0.5 1.0

The Problem

$$\nabla_s^2 f(\varphi, \vartheta) = \lambda f(\varphi, \vartheta)$$

Current solution

For $m \in \mathbb{Z}$ and $m < n$:

$$\tilde{Y}_{m,n}(\varphi, \vartheta) = \Phi(\varphi)\Theta(\vartheta) = e^{im\varphi}P_{m,n}(\cos \vartheta)$$

Intuition of conditions for m and n

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Basis functions?

The functions $\tilde{Y}_{m,n}$ span the space of nice functions $S^2 \rightarrow \mathbb{C}$.

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Definition

The inner product of nice functions $f(\varphi, \vartheta)$ and $g(\varphi, \vartheta)$ from S^2 to \mathbb{C} is

$$\langle f, g \rangle = \iint_{S^2} f g^* d\Omega$$

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The inner product of nice functions $f(\varphi, \vartheta)$ and $g(\varphi, \vartheta)$ from S^2 to \mathbb{C} is

$$\langle f, g \rangle = \iint_{S^2} f g^* d\Omega = \int_0^{2\pi} \int_0^{\pi} f(\varphi, \vartheta) g^*(\varphi, \vartheta) \sin \vartheta d\vartheta d\varphi$$

Definition

A set of basis functions are *orthonormal* if

$$\langle B_{m,n}, B_{m',n'} \rangle = \begin{cases} 1 & m = m' \wedge n = n' \\ 0 & \text{else} \end{cases}$$

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Problem

$$\langle \tilde{Y}_{m,n}, \tilde{Y}_{m',n'} \rangle = \begin{cases} \frac{4\pi}{2n+1} \frac{(n+m)!}{(n-m)!} & m = m' \wedge n = n' \\ 0 & \text{else} \end{cases}$$

Spherical Harmonics

Definition

The orthonormal spherical harmonics are

$$Y_{m,n}(\varphi, \vartheta) = N_{m,n} e^{im\varphi} P_{m,n}(\cos \vartheta)$$

where the normalisation constant

$$N_{m,n} = \sqrt{\frac{2n+1}{4\pi} \frac{(n-m)!}{(n+m)!}}$$

Fixed

$$\langle Y_{m,n}, Y_{m',n'} \rangle = \begin{cases} 1 & m = m' \wedge n = n' \\ 0 & \text{else} \end{cases}$$

Theorem

For nice periodic functions on S^2 :

$$f(\varphi, \vartheta) = \sum_{m \in \mathbb{Z}} \sum_{n \in \mathbb{Z}} c_{m,n} Y_{m,n}(\varphi, \vartheta)$$

where

$$c_{m,n} = \langle f, Y_{m,n} \rangle.$$

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Linear and Rotational Kinetic Energy

Momentum and KE

$$\mathbf{p} = m\mathbf{v}, \quad E_k = \frac{\mathbf{p}^2}{2m}$$

Linear and Rotational Kinetic Energy

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Angular M. and KE

$$\mathbf{L} = \mathbf{r} \times \mathbf{p}, \quad E_{k,a} = \frac{\mathbf{L}^2}{2mr^2}$$

Linear and Rotational Kinetic Energy

Momentum and KE

$$\mathbf{p} = m\mathbf{v}, \quad E_k = \frac{\mathbf{p}^2}{2m}$$

QM Formulation

$$\hat{\mathbf{p}} = -i\hbar\nabla, \quad \hat{E}_k = -\frac{\hbar^2}{2m}\nabla^2$$

Angular M. and KE

$$\mathbf{L} = \mathbf{r} \times \mathbf{p}, \quad E_{k,a} = \frac{\mathbf{L}^2}{2mr^2}$$

Linear and Rotational Kinetic Energy

Momentum and KE

$$\mathbf{p} = m\mathbf{v}, \quad E_k = \frac{\mathbf{p}^2}{2m}$$

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Angular M. and KE

$$\mathbf{L} = \mathbf{r} \times \mathbf{p}, \quad E_{k,a} = \frac{\mathbf{L}^2}{2mr^2}$$

QM Formulation

Pretty long derivation yields:

$$\hat{E}_{k,a} = -\frac{\hbar^2}{2mr^2}\nabla_s^2$$

Intuition for the Operators

Time independent SE

$$\hat{\mathcal{H}}|\Psi\rangle = E|\Psi\rangle$$

Time independent SE

$$\left(\hat{E}_k + U\right) |\Psi\rangle = E|\Psi\rangle$$

Time independent SE

$$\left(\frac{\hat{\mathbf{p}}^2}{2m} + U \right) |\Psi\rangle = E|\Psi\rangle$$

Time independent SE

$$\text{Meili} \quad \left[-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + U(x) \right] \Psi(x) = E\Psi(x)$$

Time independent SE

$$3D \quad \left[-\frac{\hbar^2}{2m} \nabla^2 + U(\mathbf{x}) \right] \Psi(\mathbf{x}) = E\Psi(\mathbf{x})$$

Time independent SE

$$\left\{ -\frac{\hbar^2}{2m} \frac{1}{r^2} \left[\nabla_s^2 - \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) \right] + U(\mathbf{r}) \right\} \Psi(\mathbf{r}) = E\Psi(\mathbf{r})$$

Time independent SE

$$\left[\frac{\hat{\mathbf{L}}^2}{2mr^2} + \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + U(\mathbf{r}) \right] \Psi(\mathbf{r}) = E\Psi(\mathbf{r})$$

Time independent SE

$$\left[\hat{E}_{k,a} + \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + U(\mathbf{r}) \right] \Psi(\mathbf{r}) = E \Psi(\mathbf{r})$$

Schrödinger Equation

Time independent SE

$$\left[\underbrace{\hat{E}_{k,a} + \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right)}_{\text{Kinetic Energy}} + U(\mathbf{r}) \right] \Psi(\mathbf{r}) = E \Psi(\mathbf{r})$$

Schrödinger Equation

Time independent SE

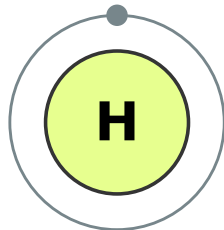
$$\left[\hat{E}_{k,a} + \underbrace{\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right)}_{\text{Radial KE } \hat{E}_{k,r}} + U(\mathbf{r}) \right] \Psi(\mathbf{r}) = E \Psi(\mathbf{r})$$

Schrödinger Equation

Time independent SE

$$\left\{ \hat{E}_{k,a} + \hat{E}_{k,r} + U(\mathbf{r}) \right\} \Psi(\mathbf{r}) = E\Psi(\mathbf{r})$$

But why?



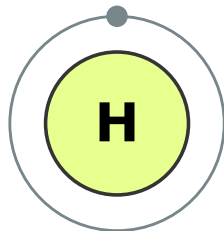
Schrödinger Equation

Time independent SE

$$\left\{ \hat{E}_{k,a} + \hat{E}_{k,r} + U(\mathbf{r}) \right\} \Psi(\mathbf{r}) = E\Psi(\mathbf{r})$$

But why?

Hydrogen atom has radial symmetry!



Electron Orbitals

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